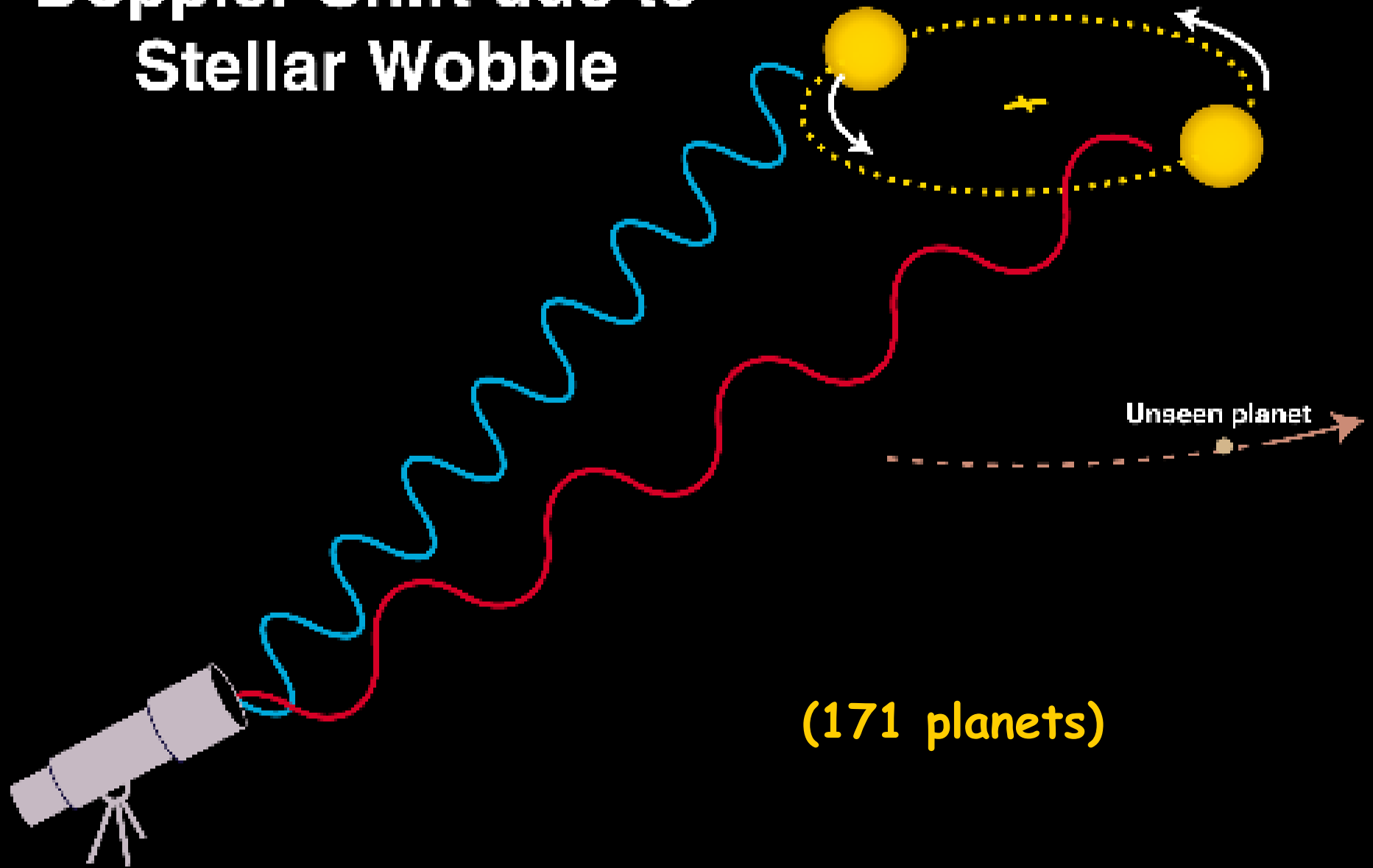


## **2. Extrasolar planets**

**Frontiers of Astronomy Workshop/School  
Bibliotheca Alexandrina  
March-April 2006**

# Doppler Shift due to Stellar Wobble



(171 planets)

For a planet of mass  $m$  on a circular orbit of semi-major axis  $a$  around a star of mass  $M \gg m$ ,

$$v_{\text{planet}} = \sqrt{\frac{GM}{a}} \sin \frac{2\pi(t-t_0)}{P}, \quad v_{\text{star}} = -\frac{m}{M} v_{\text{planet}}$$

**prove!**

where  $P = 2\pi (a^3/GM)^{1/2}$ . This neglects

- inclination (assumes orbit is edge-on)
- eccentricity

Radial velocity is

$$v_{\text{star}} = \frac{m \sin I}{M} \sqrt{\frac{GM}{a}} (1-e^2)^{-1/2} [\cos(f+\omega) + e \cos \omega]$$

where

$$\tan^{-1} \frac{1}{2} f = \sqrt{\frac{1+e}{1-e}} \tan^{-1} \frac{1}{2} u, \quad u - e \sin u = \frac{2\pi}{P} (t - t_0).$$

$$P = \sqrt{\frac{GM}{a^3}}$$

**m** = planet mass

**M** = star mass

**I** = inclination of orbit

**a** = semi-major axis

**e** = eccentricity

**$\omega$**  = argument of pericenter

**f** = true anomaly

**u** = eccentric anomaly

Radial velocity is

$$v = \frac{m \sin I}{M} \sqrt{\frac{GM}{a}} (1-e^2)^{-1/2} [\cos(f+\omega) + e \cos \omega]$$

where

$$\tan^{-1} \frac{1}{2} f = \sqrt{\frac{1+e}{1-e}} \tan^{-1} \frac{1}{2} u, \quad u - e \sin u = \frac{2\pi}{P} (t - t_0).$$

$$P = \sqrt{\frac{GM}{a^3}}$$

Given the star mass  $M$  (known from spectral type), radial-velocity observations yield  $a$ ,  $m \sin I$ ,  $e$ ,  $\omega$

# Observational techniques

- spectrograph with resolving power of 100,000 has pixel scale 3 km/s. Best observers can now measure  $\Delta v \sim 3$  m/s or 1/1000 of a pixel and are stable over many years.
  - use many lines in spectrum
  - very high signal-to-noise (200-500)
  - pass light through iodine cell, so calibration lines are distorted in the same way as the data
- error sources:
  - photon noise (use largest telescopes, bright stars)
  - weak or broad spectral lines (focus on old G stars)
  - stellar activity (convection, p-modes)

### 51 Peg:

$$m \sin I = 0.46 M_J$$

$$P = 4.23 \text{ d}$$

$$a = 0.05 \text{ AU}$$

$$e = 0.01$$

### 70 Vir:

$$m \sin I = 7.44 M_J$$

$$P = 116.7 \text{ d}$$

$$a = 0.48 \text{ AU}$$

$$e = 0.40$$

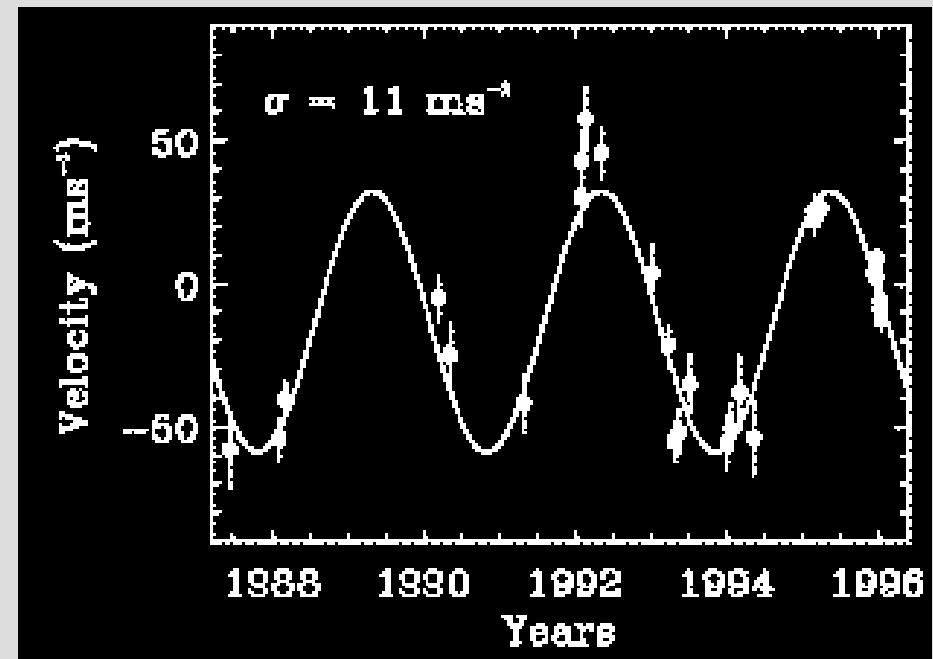
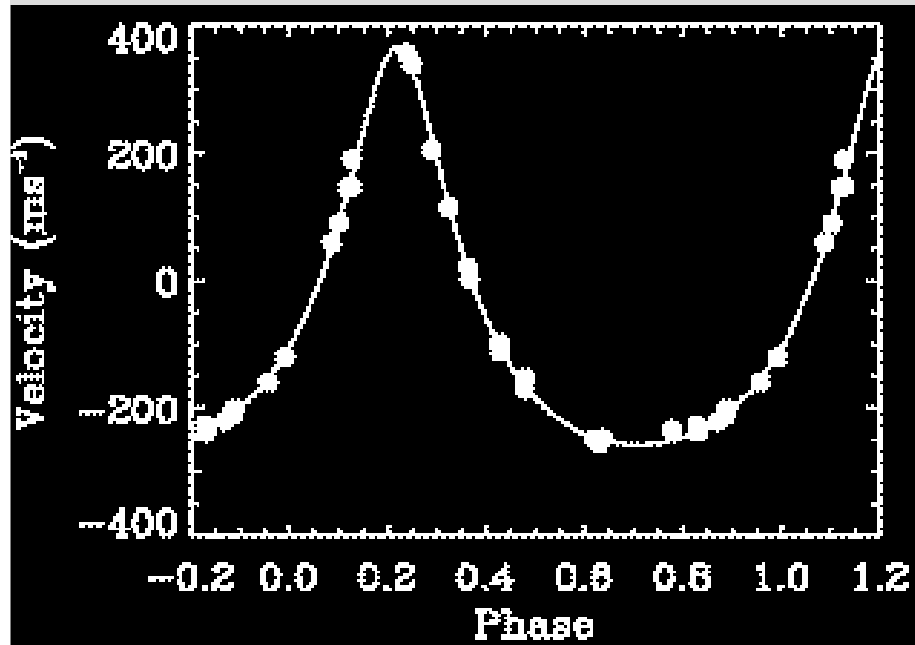
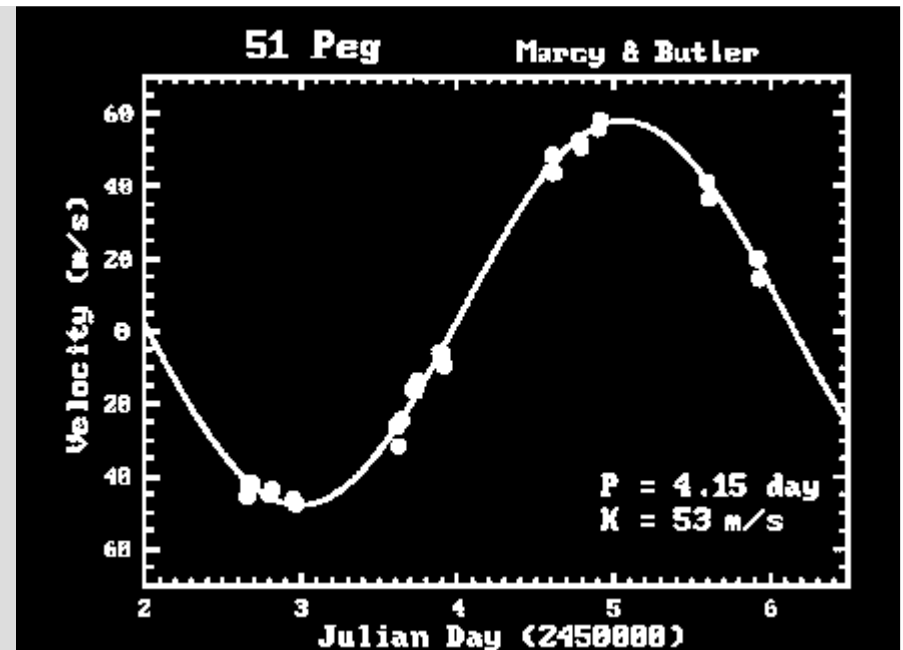
### 47 UMa:

$$m \sin I = 2.54 M_J$$

$$P = 1089 \text{ d}$$

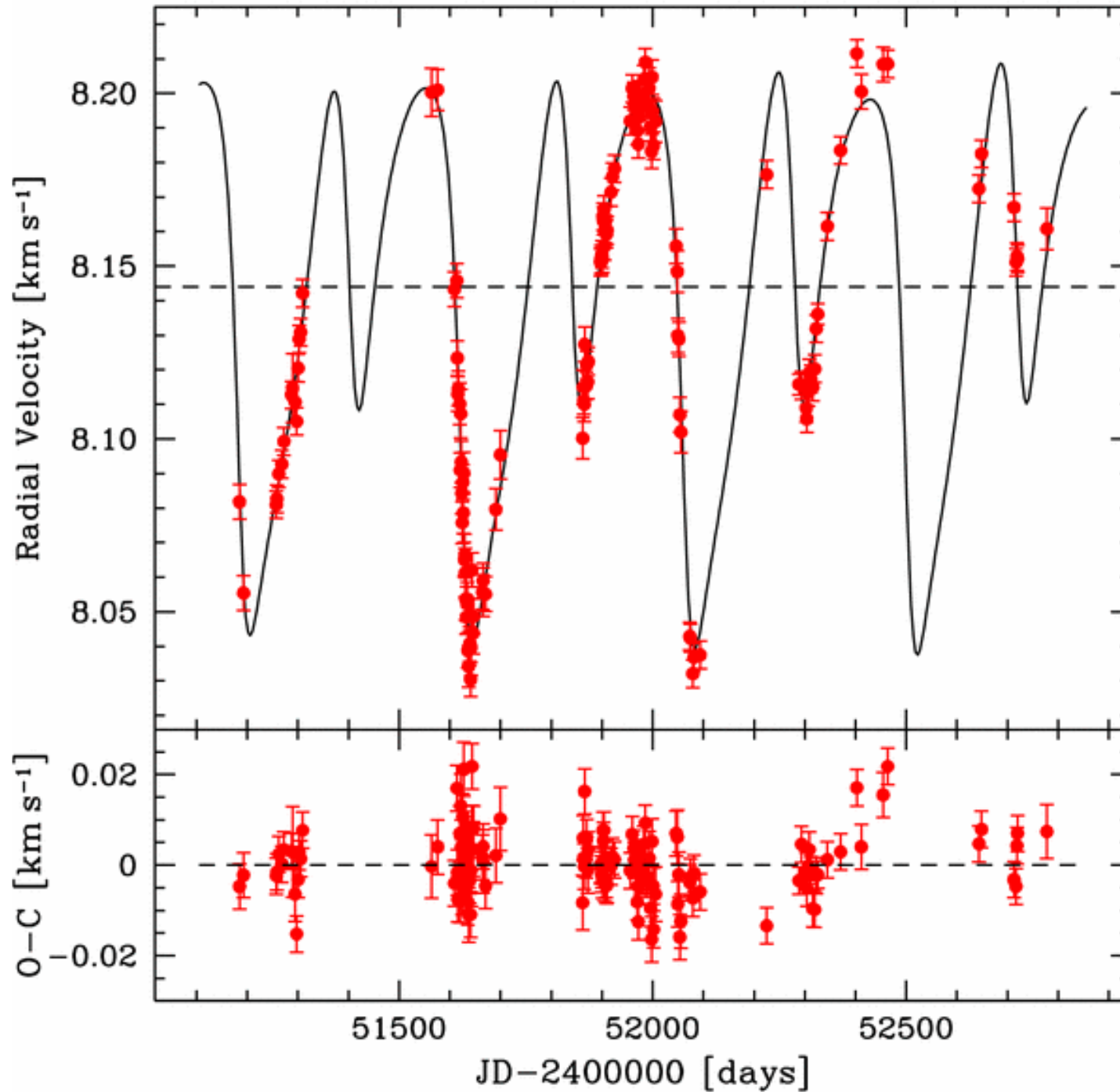
$$a = 2.09 \text{ AU}$$

$$e = 0.06$$



HD 82943

CORALIE



HD 82943

planet 1:

 $m \sin I = 1.84 m_J$  $P = 435 \text{ d}$  $e = 0.18 \pm 0.04$ 

planet 2:

 $m \sin I = 1.85 m_J$  $P = 219 \text{ d}$  $e = 0.38 \pm 0.01$ 

(Mayor et al. 2003)



# Pulsar planets

Arrival time is related to emission time by

$$t = t_e + cr$$

where  $r$  is the distance to the pulsar. Observed pulsar period is

$$\Delta t = \Delta t_e + \Delta r/c = \Delta t_e + v_r \Delta t_e/c = \Delta t_e(1 + v_r/c)$$

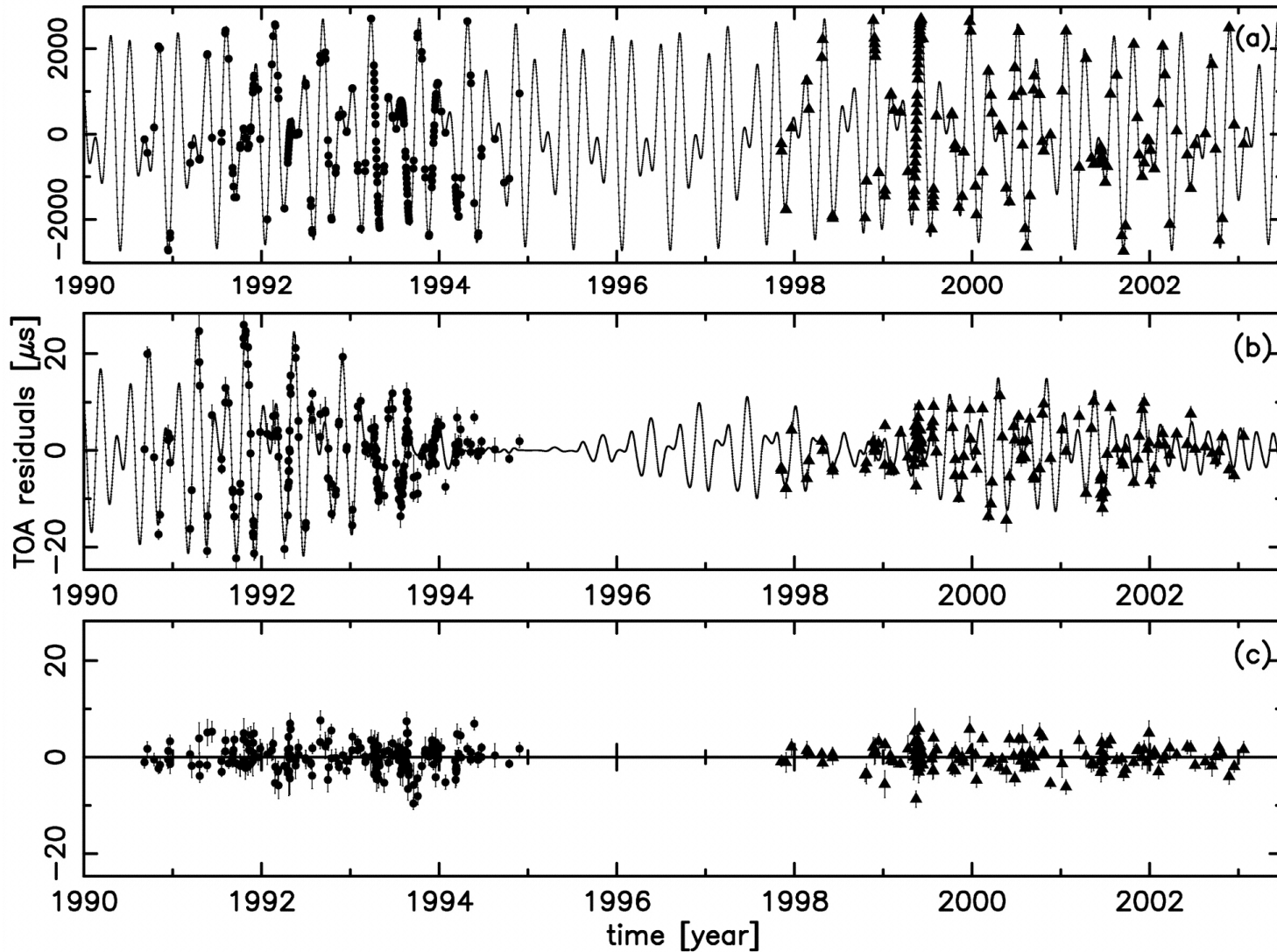
where  $\Delta t_e$  is the period of the pulsar in its rest frame. Rate of change of observed period is

$$\frac{d\Delta t}{dt} = \frac{\Delta t_e}{c} \frac{dv_r}{dt}$$

# Pulsar planets

- three planets discovered orbiting PSR B1257+12 by Wolszczan & Frail (1992)
- orbital parameters can be determined far more accurately than for radial-velocity measurements of nearby stars
- two planets near 3:2 resonance which enhances mutual perturbations, so these can be measured
- remarkably similar to the inner solar system

PSR B1257+12, Arecibo, 430 MHz



- (a) no planets
- (b) three planets
- (c) three planets + mutual interactions

Konacki & Wolszczan (2003)

TABLE 2  
ORBITAL AND PHYSICAL PARAMETERS OF PLANETS<sup>a</sup>

Parameter	Planet A	Planet B	Planet C
Projected semimajor axis, $x^0$ (ms) .....	0.0030 (1)	1.3106 (1)	1.4134 (2)
Eccentricity, $e^0$ .....	0.0	0.0186 (2)	0.0252 (2)
Epoch of pericenter, $T_p^0$ (MJD) .....	49765.1 (2)	49768.1 (1)	49766.5 (1)
Orbital period, $P_b^0$ (day) .....	25.262 (3)	66.5419 (1)	98.2114 (2)
Longitude of pericenter, $\omega^0$ (deg) .....	0.0	250.4 (6)	108.3 (5)
Mass ( $M_\oplus$ ) .....	0.020 (2)	4.3 (2)	3.9 (2)
Inclination, solution 1, $i^0$ (deg) .....	...	53 (4)	47 (3)
Inclination, solution 2, $i^0$ (deg) .....	...	127 (4)	133 (3)
Planet semimajor axis, $a_p^0$ (AU) .....	0.19	0.36	0.46
Non-Keplerian dynamical parameters .....	...	...	...
$\gamma_B$ ( $\times 10^{-6}$ ) .....	...	9.2 (4)	...
$\gamma_C$ ( $\times 10^{-6}$ ) .....	...	8.3 (4)	...
$\tau$ (deg) .....	...	2.1 (9)	...

<sup>a</sup> Figures in parentheses are the formal  $1 \sigma$  uncertainties in the last digits quoted.

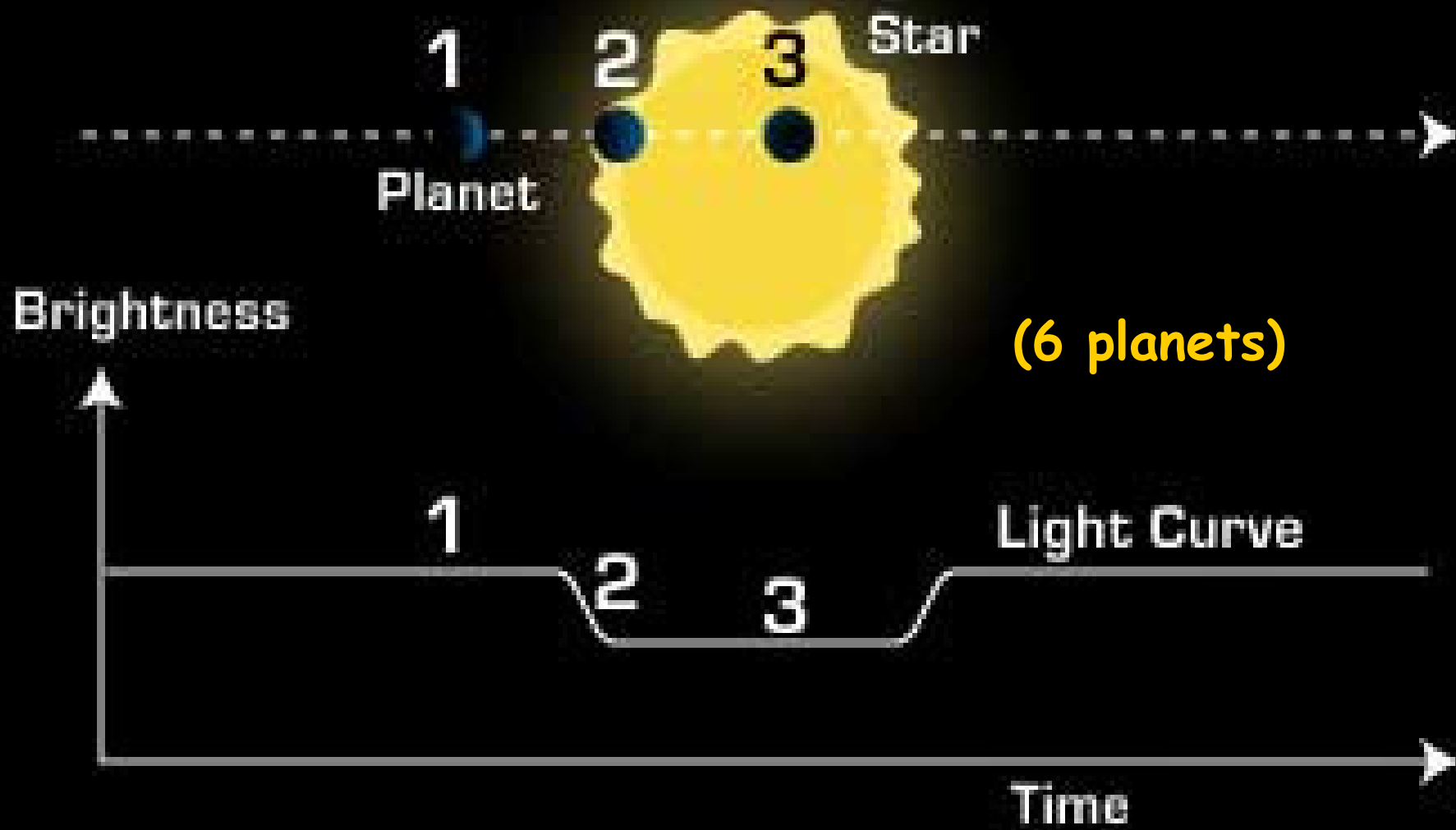
**Konacki & Wolszczan (2003)**

## Radial-velocity surveys - summary

$$v \propto \frac{m \sin I}{M} \sqrt{\frac{GM}{a}}$$

- sensitive to high-mass planets, small semi-major axes, short orbital periods
- give only  $m \sin I$ , not  $m$
- maximum sensitivity  $\Delta v \sim 2 \text{ m/s}$
- Jupiter  $\Delta v \sim 13 \text{ m/s}$ ,  $P = 11.9 \text{ yr}$  - just detectable
- Earth  $\Delta v \sim 0.1 \text{ m/s}$ ,  $P = 1 \text{ yr}$  - **not** detectable

# Transit Method



# Transit surveys

## Primary eclipse:

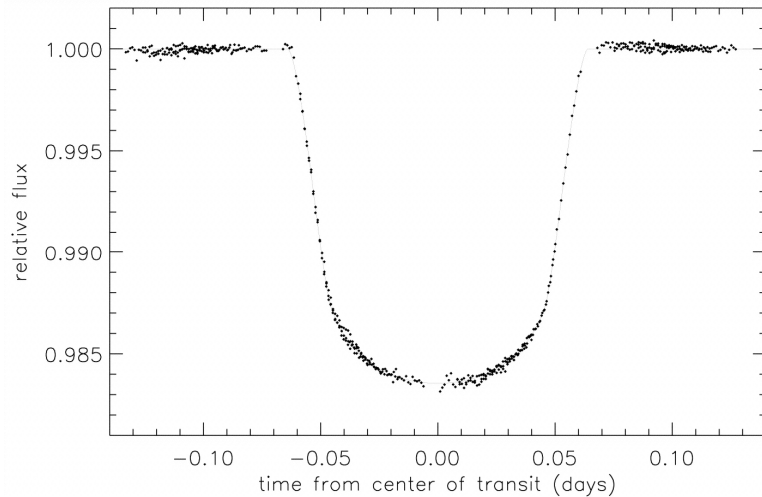
- planet passes **in front of** the star **prove!**
- flux from star is reduced by  $(R_{\text{planet}}/R_{\text{star}})^2$ . For a solar-type star this is 1% for Jupiter or 0.01% for Earth

## Secondary eclipse:

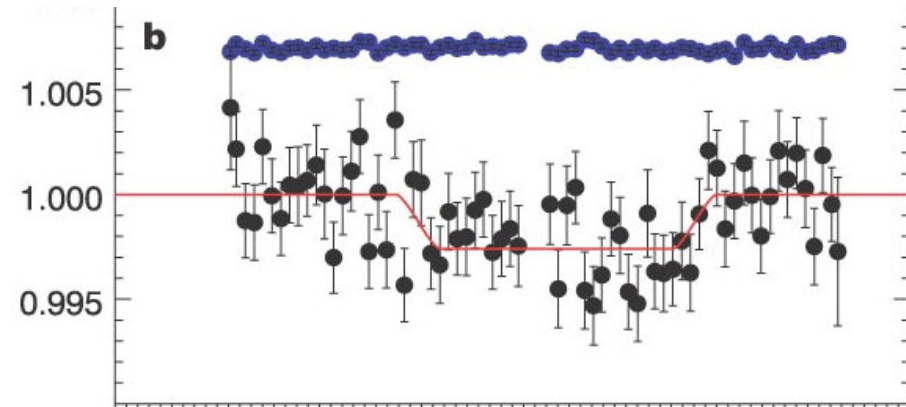
- planet passes **behind** the star and flux is decreased by eclipse of light from planet
- in visible light, planet emits by reflection from surface: total flux is reduced by  $p(R_{\text{planet}}/a)^2$  where  $p$  is albedo (0.4-0.6 for giant planets in solar system) **prove!**
- in infrared, planet emits thermally as black body. If planet and star are black bodies in Rayleigh-Jeans limit ( $\lambda > 3\mu(1000 \text{ K}/T)$ ) then flux reduced by  $(R_{\text{planet}}/R_{\text{star}})^2(T_{\text{planet}}/T_{\text{star}})^2$

**prove!**

## primary (visible)



## secondary (infrared)



- mass = 0.69 Jupiter masses
- radius = 1.35 Jupiter radii (“bloated”)
- orbital period 3.52 days, orbital radius 0.047 AU or 10 stellar radii
- stellar obliquity  $< 10^\circ$
- $T = 1130 \pm 150$  K
- sodium, oxygen, carbon detected from planetary atmosphere

HD 209458

Brown et al. (2001),  
Deming et al. (2005)



# Transit searches

## Ground-based:

ASP, BEST, GITPO,  
HATNetwork,  
MONET, OGLE, PASS,  
PISCES, STARE,  
STEPSS, TAPT, TEP,  
Transitsearch.org,  
UstAPS, Vulcain,  
Vulcain South, WHAT,  
STELLA, SuperWASP

## Space missions:

COROT (2006), GAIA  
(2011), GEST, Kepler  
(2008)

current record: 5  
from OGLE, 1 from  
STARE, 3 from  
radial-velocity  
surveys

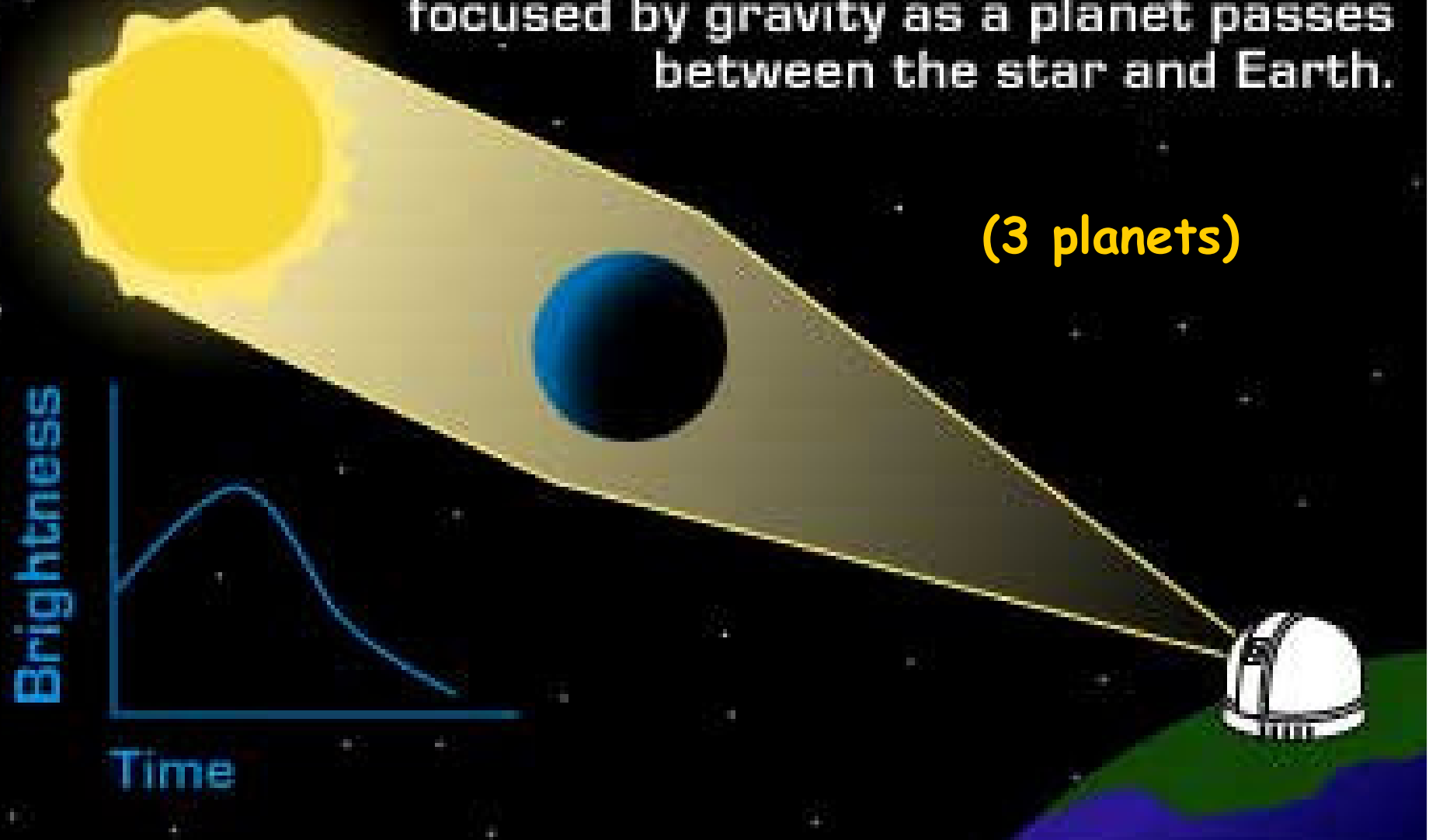
# Transit searches

Why are these so hard?

- probability that a given planet will transit is small,  $\sim R_{\text{star}}/a$  (only 0.5% at  $a=1$  AU) **prove!**
- transit duration is  $(R_{\text{star}}/a)P/\pi$  **prove!**
- transit depth is small,  $<1\%$
- confusion from grazing eclipsing binary stars
- star spots, stellar pulsations, stellar flares

# Gravitational Microlensing

Light from a distant star is bent and focused by gravity as a planet passes between the star and Earth.

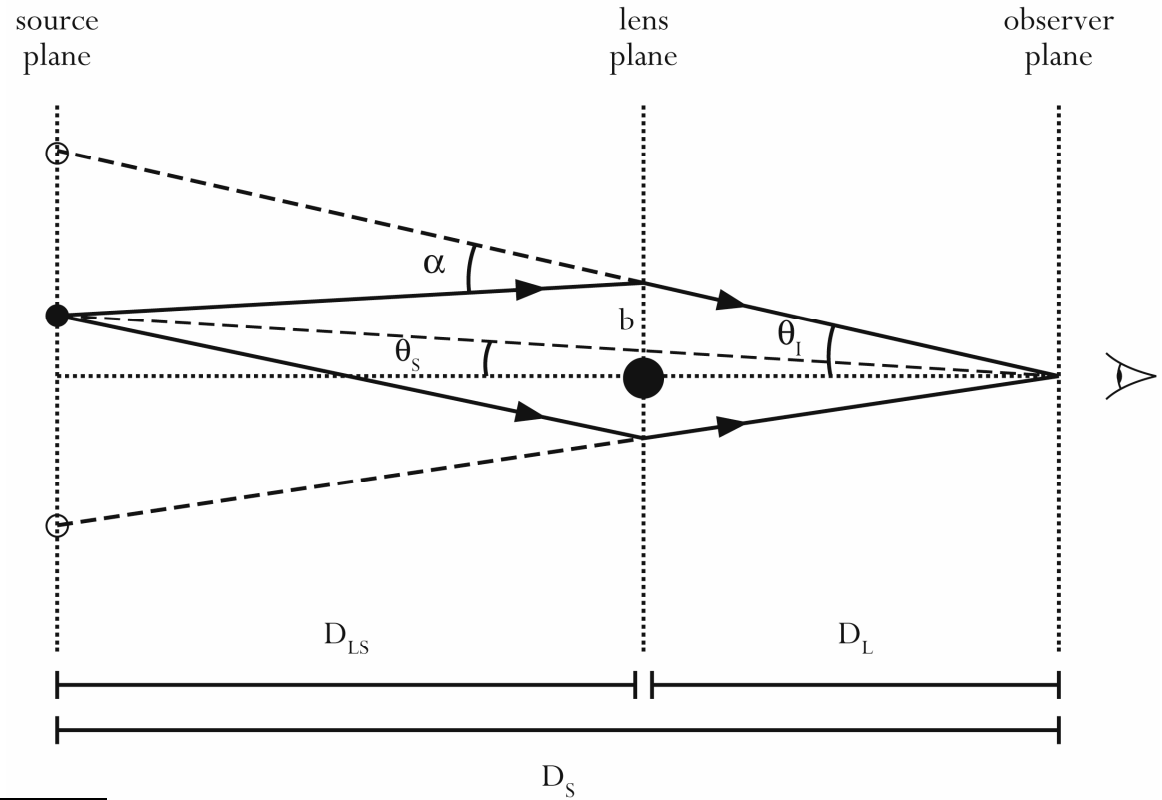


# Gravitational lensing

- first proposed by Einstein in 1936
- a particle travelling at high speed  $v$  past a mass  $M$  with impact parameter  $b$  suffers angular deflection  $\alpha = 2GM/v^2b$
- in general relativity, deflection of a photon is obtained by replacing  $v$  by  $c$  and multiplying by 2:

$$\alpha = \frac{4GM}{c^2b}$$

# Gravitational lensing



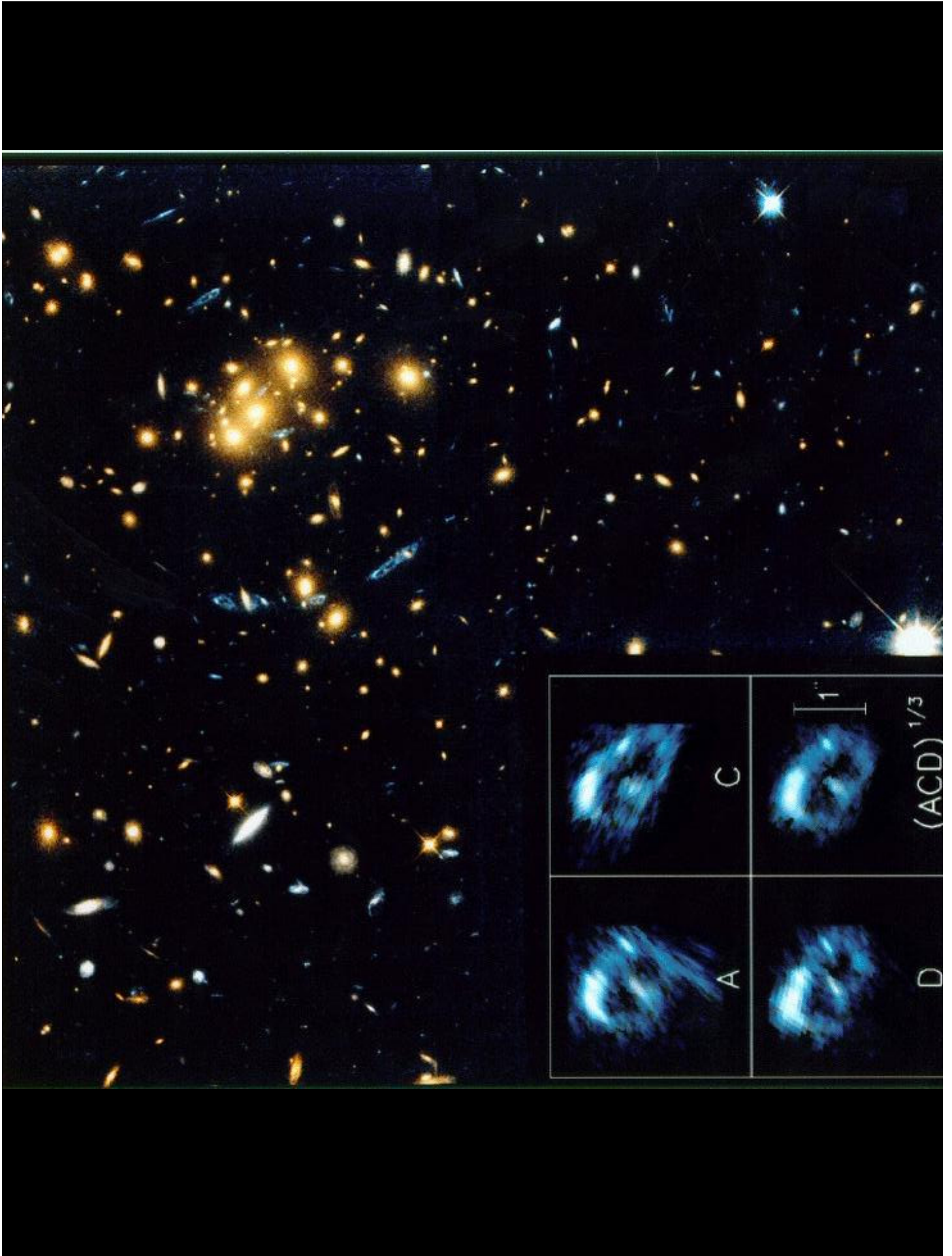
Einstein ring radius

$$\theta_E = \sqrt{\frac{4GM}{c^2} \frac{D_{SL}}{D_L D_S}}$$

Then observed angle  $\theta$  is related to true source angle  $\theta_S$  by the lensing equation

$$\theta_S = \theta - \frac{\theta_E^2}{\theta} \quad \text{prove!}$$

If  $\theta_S = 0$  (observer, lens and source are collinear) then  $\theta = \theta_E$ .



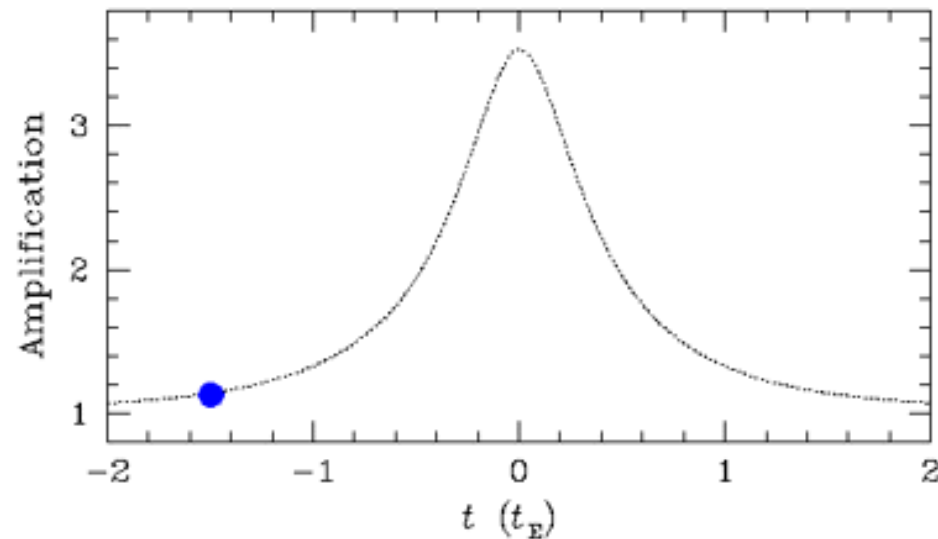
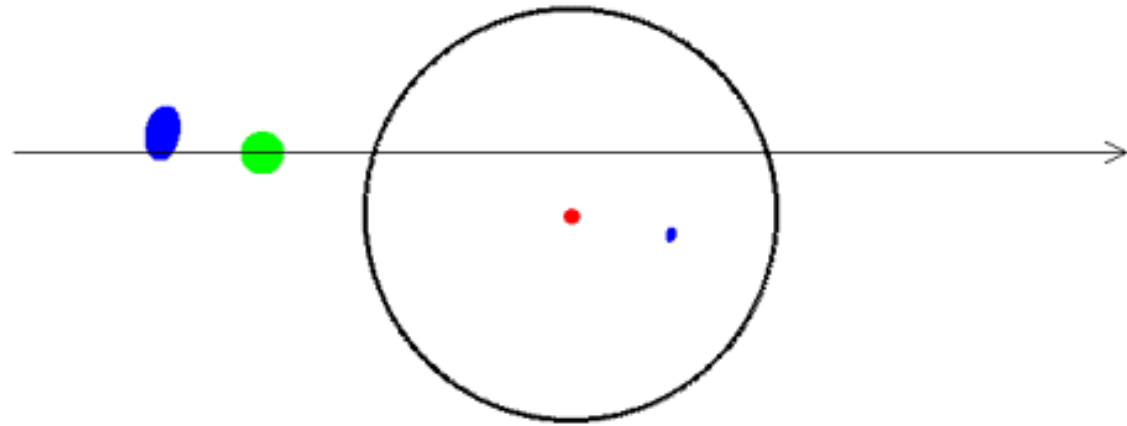
# Gravitational lensing

surface brightness is conserved so distortion of image of source across larger area of sky implies magnification

$$A = \frac{u^2 + 2}{u\sqrt{u^2 + 4}},$$

$$u \equiv \frac{\theta_S}{\theta_E}$$

$$\beta = 0.3$$
$$r_s = 0.1 \theta_E$$



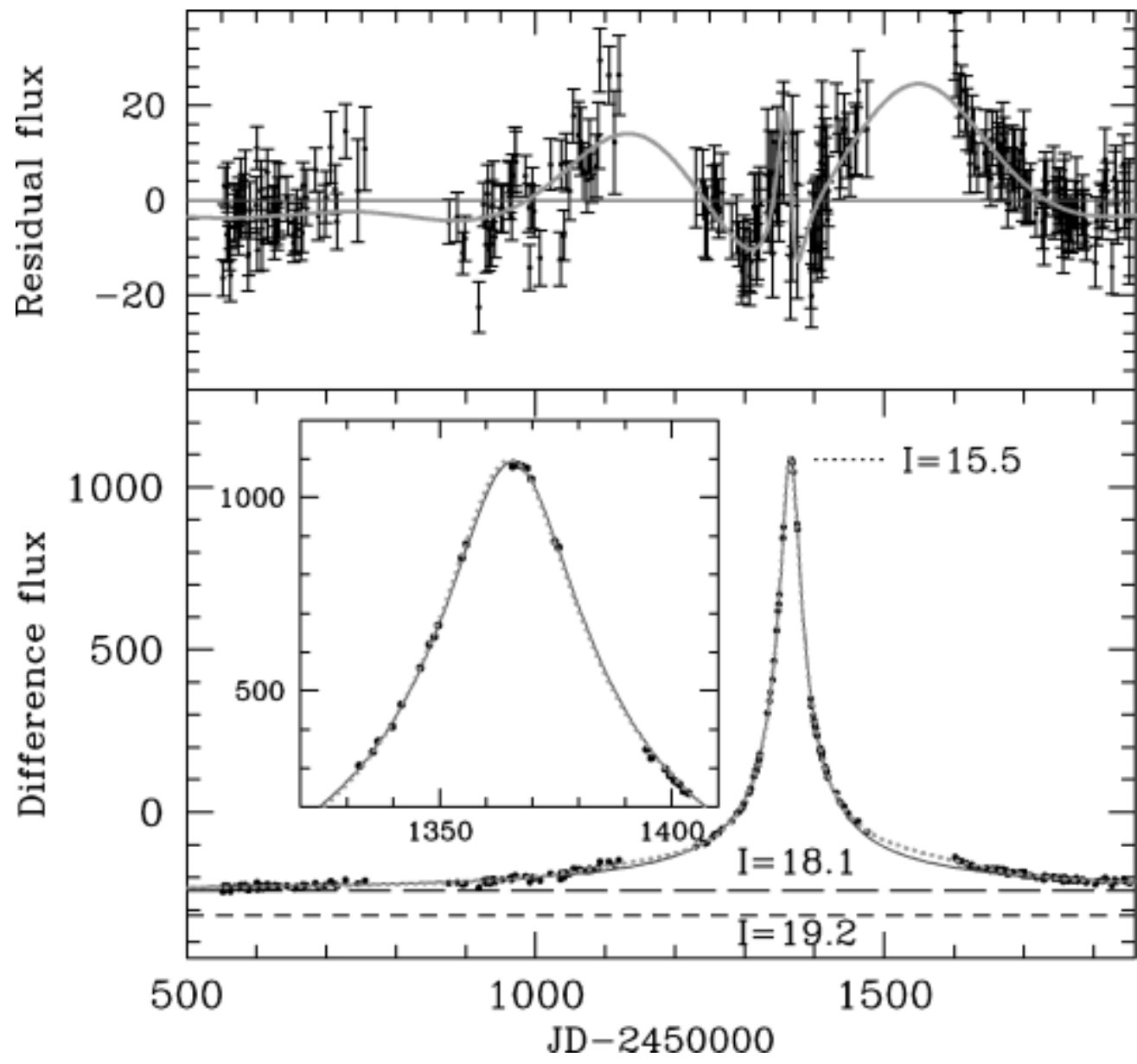
**prove!**

# Gravitational microlensing

Consider a source star in the Galactic bulge at  $D_S \sim 7$  kpc, lensed by an intervening star at  $D_L \sim 4$  kpc. For solar-type lens star,  $\theta_E = 0.001$  arcsec  $\sim 4$  AU. **prove!**

- image splitting or shift is impossible to see
- image magnification is easy to see
- time required to transit Einstein ring  $\sim D_L \theta_E / v \sim 0.2$  yr, for  $v \sim 100$  km/s **prove!**
- substantial magnification if and only if impact parameter less than Einstein radius
- chance that any given star is microlensed is only one per million

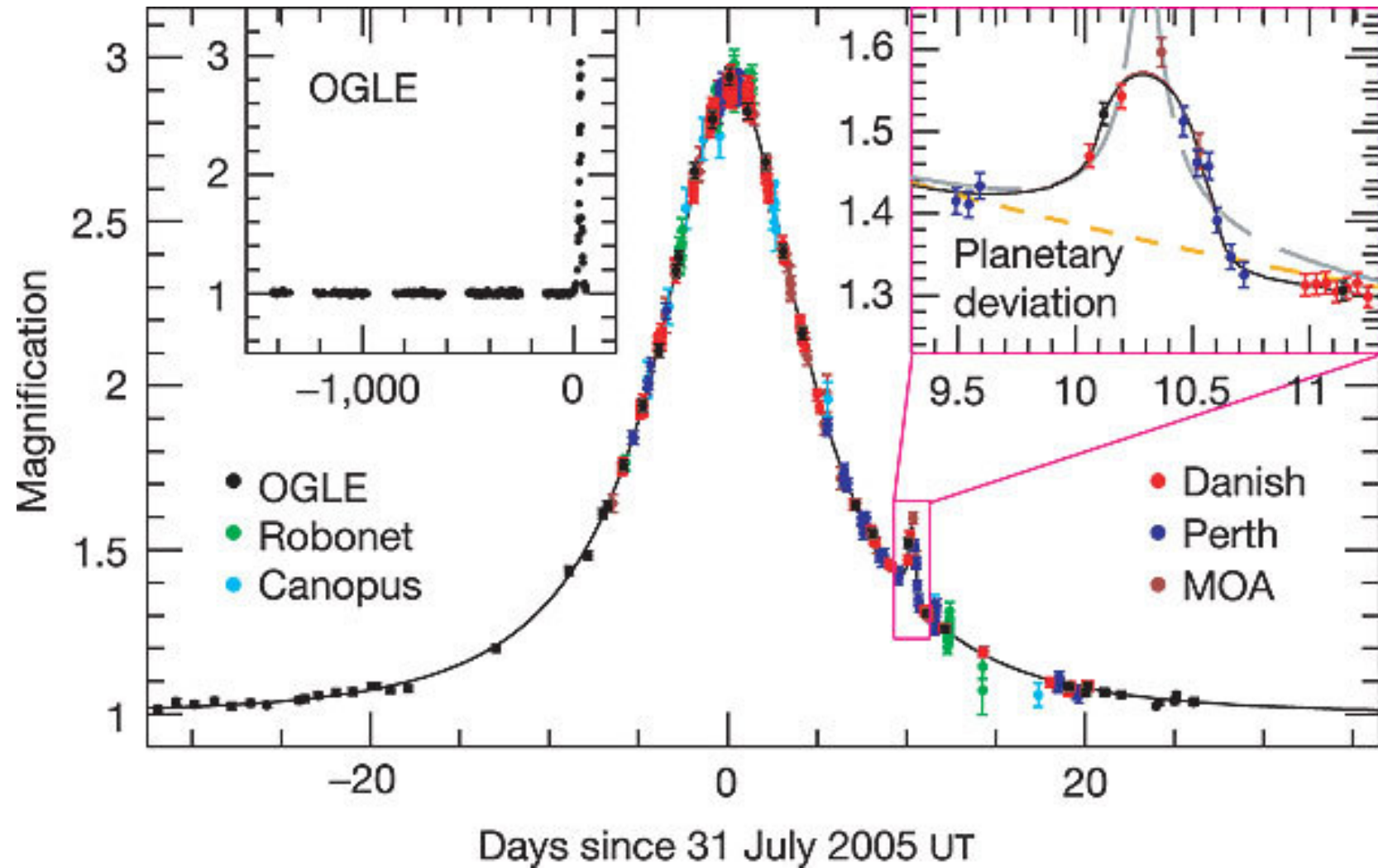




Mao et al. (2002)

# Gravitational microlensing of planets

- Einstein radius scales as  $M^{1/2}$  so cross-section and expected duration scale as  $M^{1/2} \sim 0.03$  for Jupiter, i.e. duration  $\sim 1$  day for Jupiter,  $\sim 1$  hour for Earth
- image magnification is the same
- Einstein ring radius  $\sim$  typical planet orbital radius
- use two-step strategy:
  1. wide-angle surveys find microlensed stars with coarse time resolution
  2. small telescopes monitor microlensed stars constantly

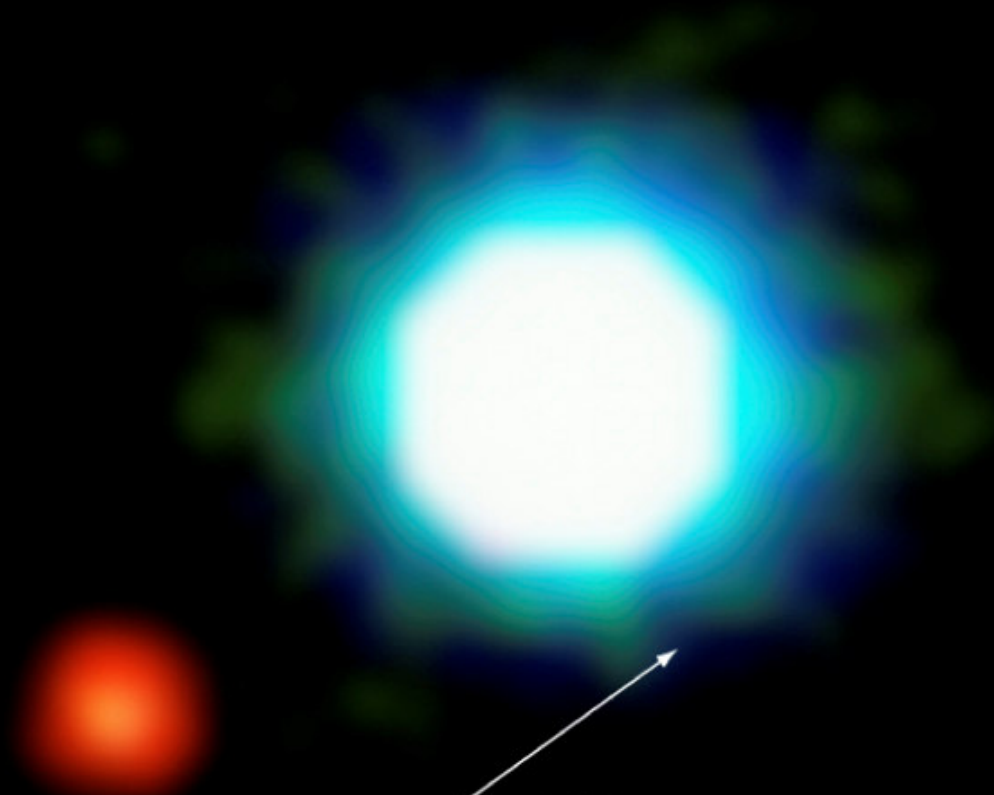


Beaulieu et al. (2006):  $5.5 (+5.5/-2.7) M_{\text{Earth}}$ ,  $2.6(+1.5/-0.6) \text{ AU}$  orbit,  $0.22(+0.21/-0.11) M_{\text{Sun}}$ ,  $D_L = 6.6 \pm 1.1 \text{ kpc}$

2MASSWJ1207334-393254

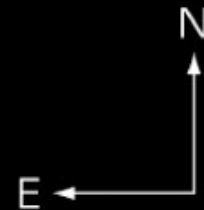
# Imaging

(1 planet)



778 mas  
55 AU at 70 pc

Chauvin et al. (2005)



# What have we learned?

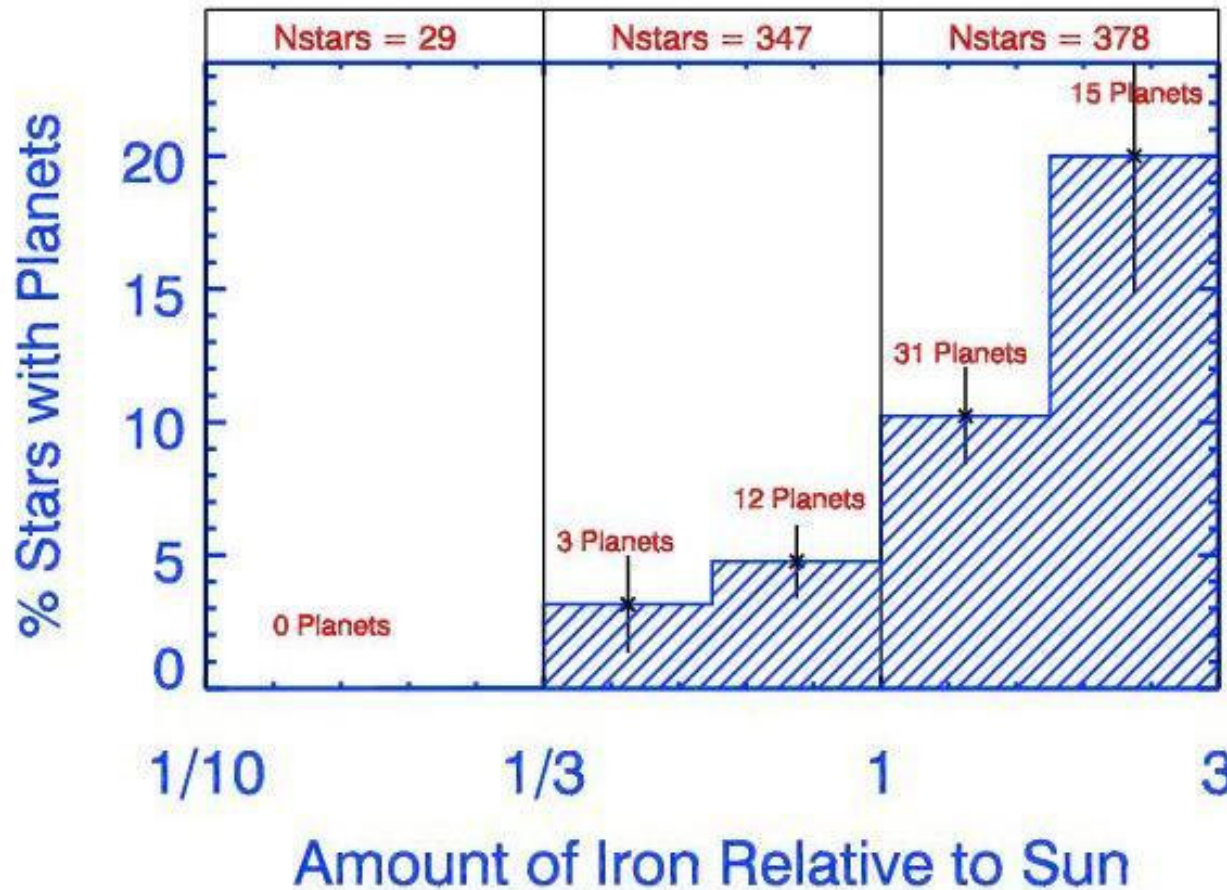
- 185 extrasolar planets known (March 2006)
  - 171 from radial-velocity surveys
  - 6 from transits
  - 4 around pulsars
  - 3 by microlensing
  - 1 by imaging

(see <http://vo.obspm.fr/exoplanetes/encyclo/encycl.html>)

# What have we learned?

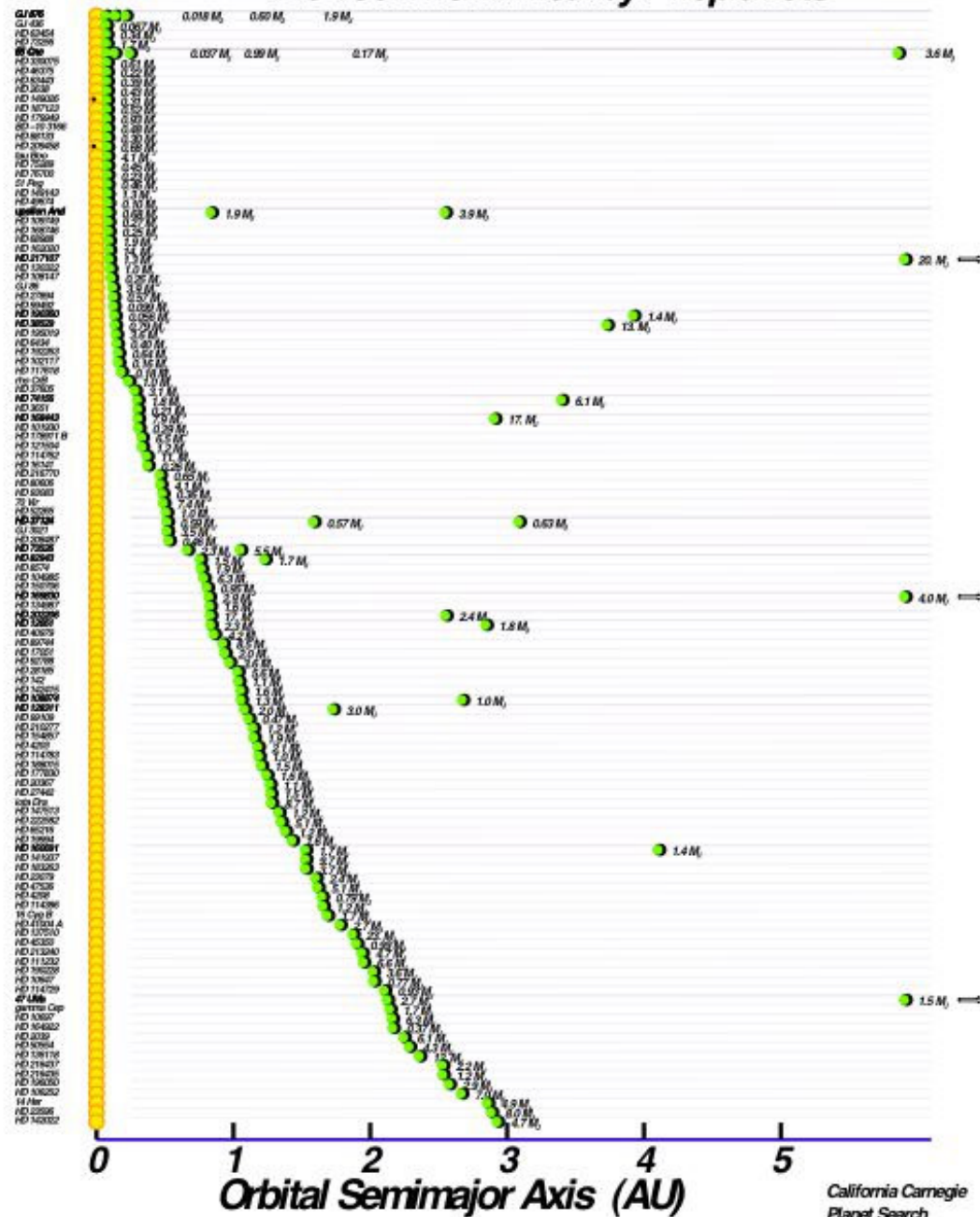
- smallest semi-major axis  $a = 0.021 \text{ AU} = 4.5 R_{\text{Sun}}$
- largest semi-major axis  $a = 5.26 \text{ AU}$  (Jupiter = 5.2 AU)
- biggest eccentricity  $e = 0.93$
- smallest eccentricity  $e = 0$
- smallest mass  $0.018 M_{\text{Jupiter}} = 5.7 M_{\text{Earth}}$
- biggest mass  $15 M_{\text{Jupiter}}$

# What have we learned?



- planets are remarkably common, especially around metal-rich stars
- probability of finding a planet  $\propto$  mass in metals in the star

## The 156 Known Nearby Exoplanets

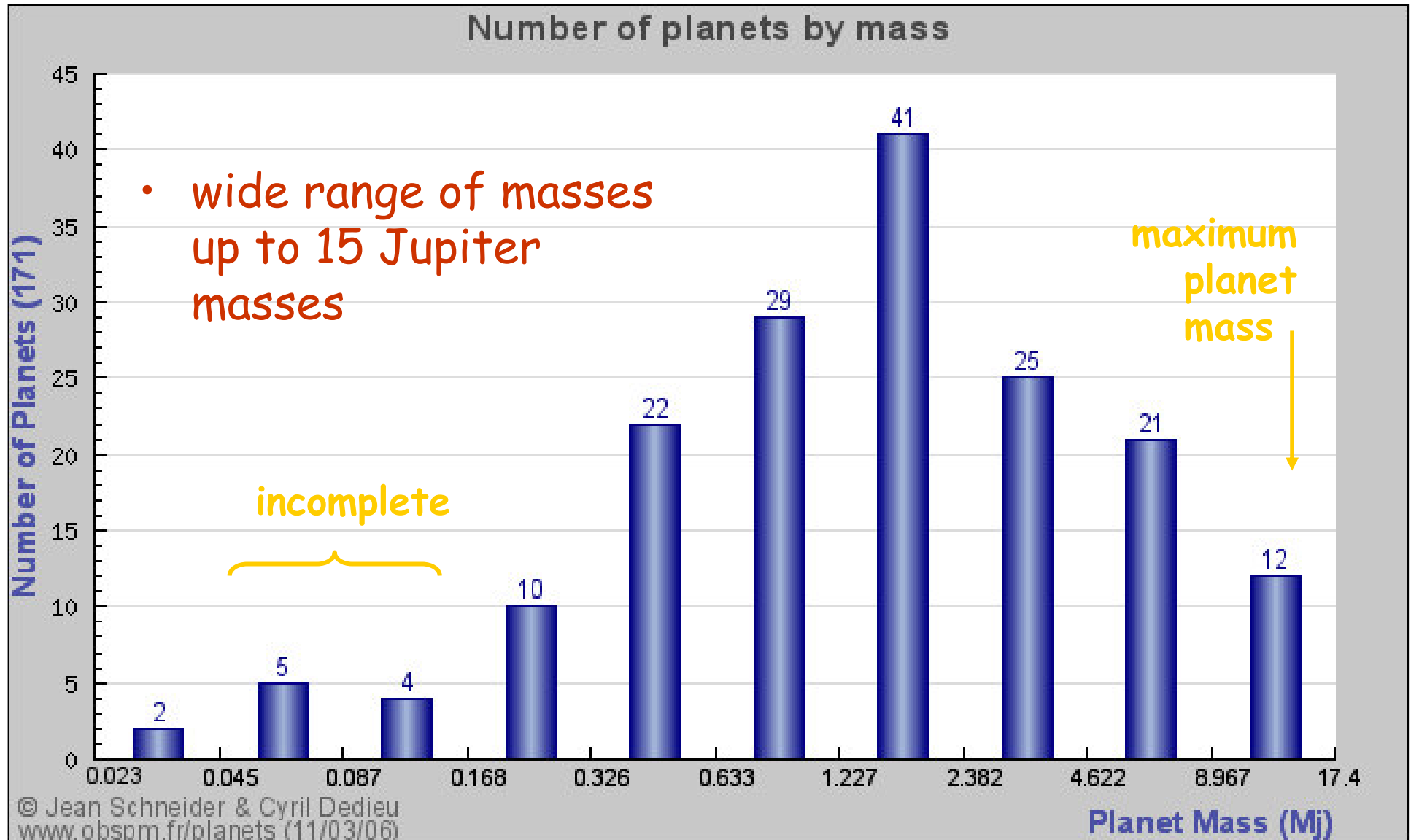


# What have we learned?

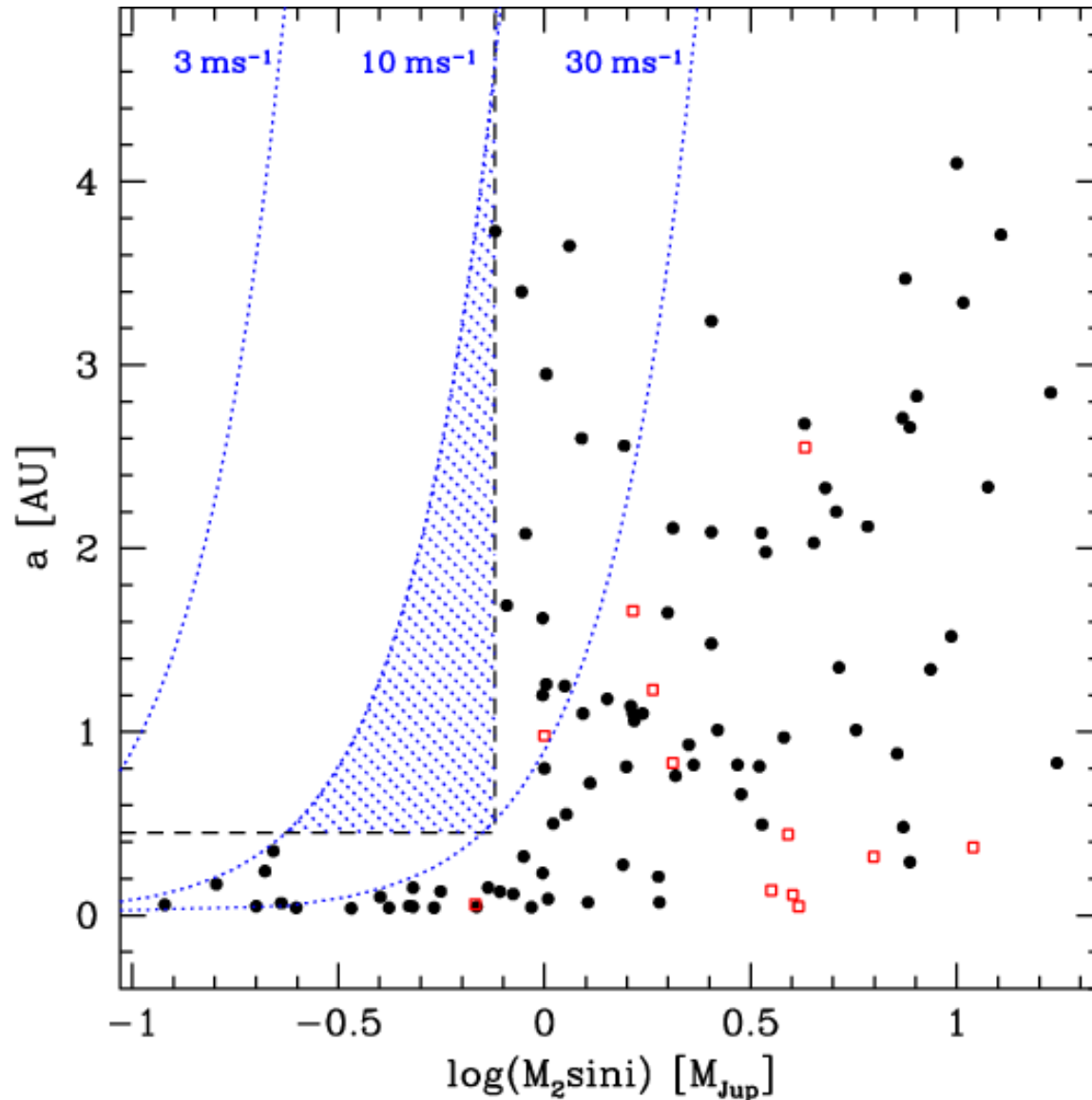
- giant planets like Jupiter and Saturn are found at *very* small orbital radii
- Gliese 876d:  $M \sin I = 0.0023 M_{\text{Jupiter}}$ ,  $P = 1.94$  days,  $a = 0.0208$  AU



# What have we learned?



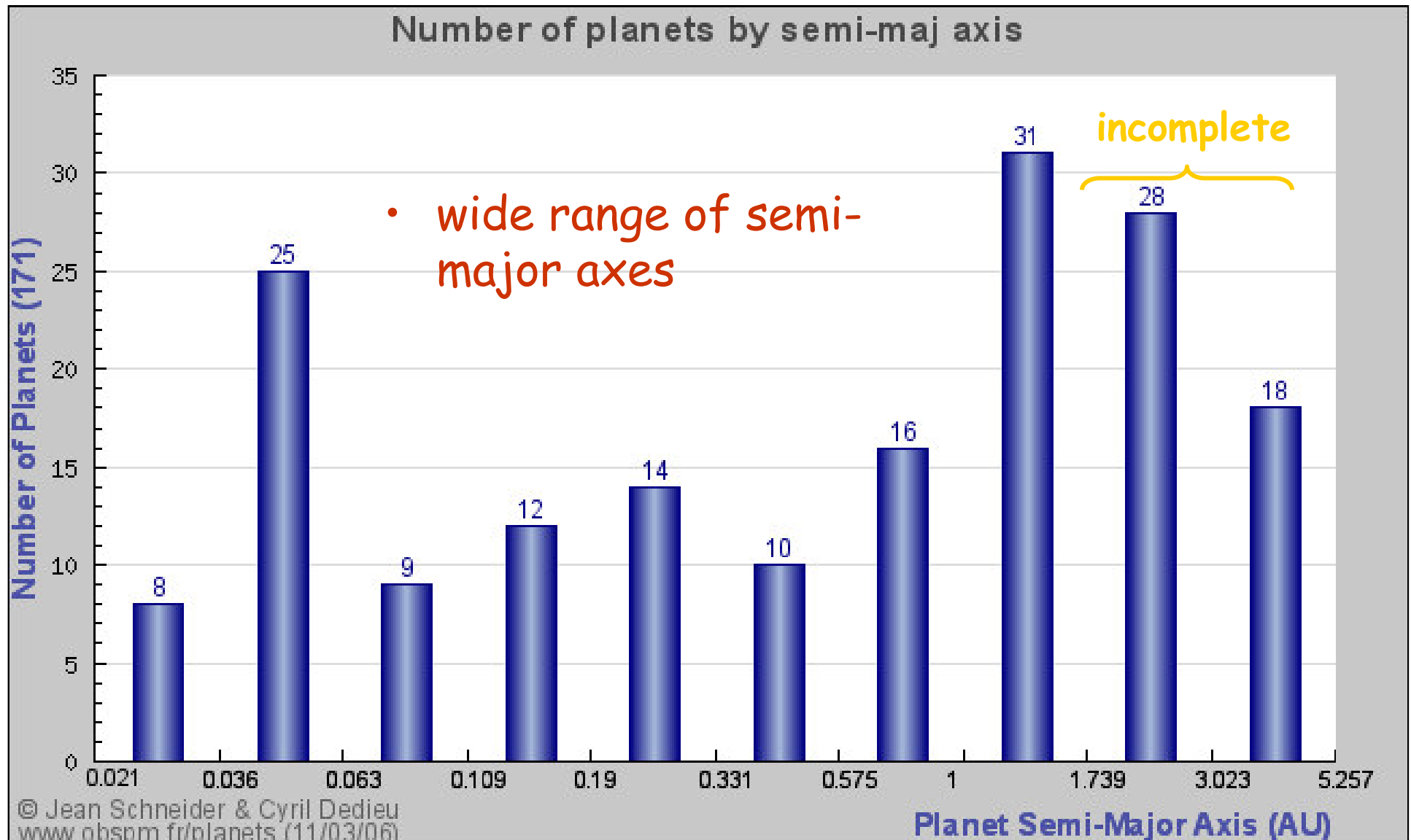
# What have we learned?



- wide range of masses up to 10 Jupiter masses

Udry et al. (2003)

# What have we learned?



# Mass and semi-major axis distribution

To a first approximation,

$$dn \propto M^{-\alpha} dM, \quad M < 10 M_J$$

$$\alpha = 1.1 \pm 0.1$$

*Planets are uniformly distributed in  $\log M$*

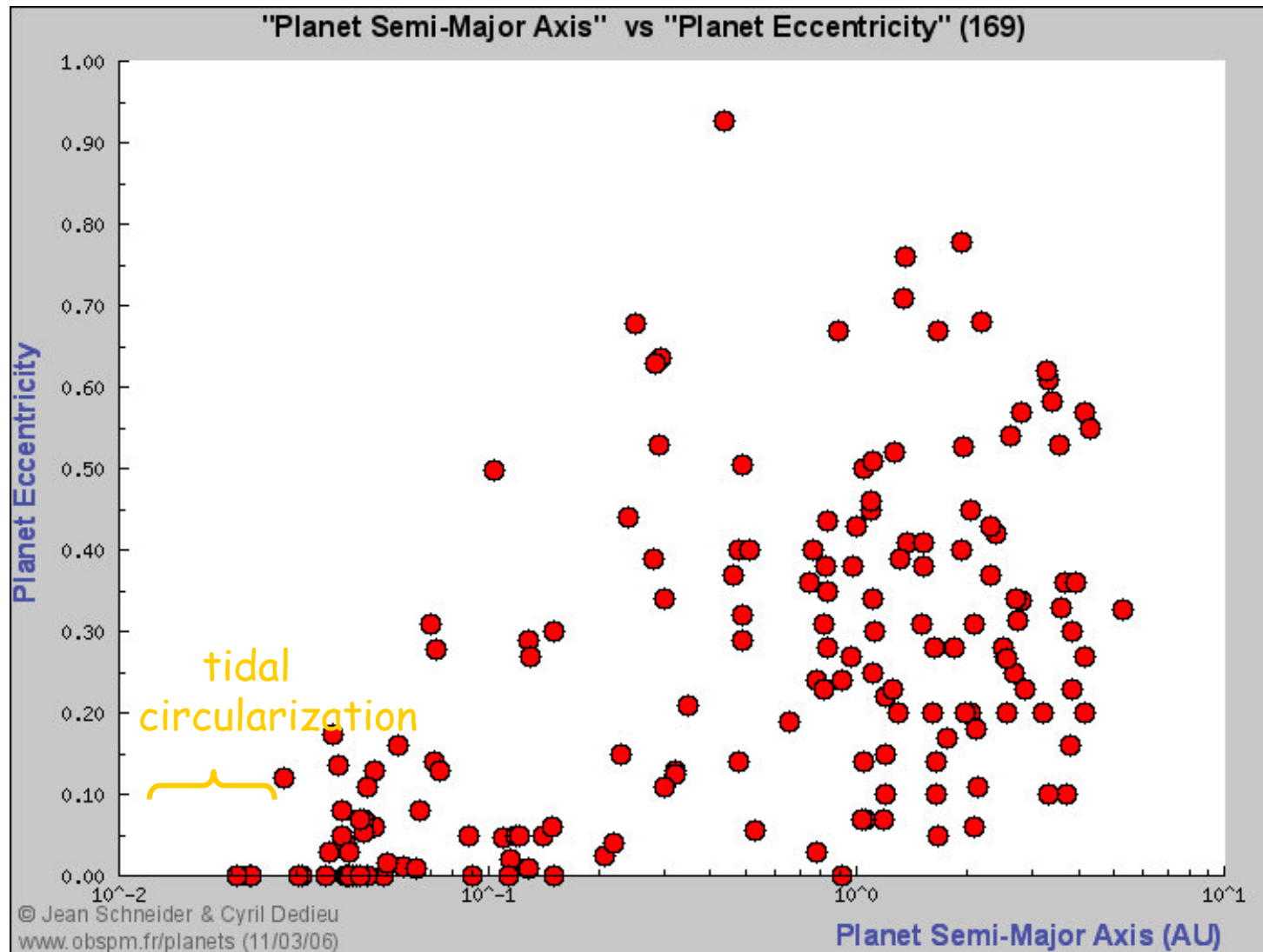
To a first approximation,

$$dn \propto a^{-\beta} da$$

$$\beta = 0.6 \pm 0.1$$

(Tabachnik & Tremaine 2002)

# What have we learned?



eccentricities  
are much  
larger than in  
the solar  
system